

# Deriving quantum mechanics from statistical assumptions

U. Klein

University of Linz, Institute for Theoretical Physics, A-4040 Linz, Austria

Joint Annual Meeting of ÖPG/SPS/ÖGAA - Innsbruck 2009

# The interpretation of quantum mechanics

This talk is about the interpretation of Quantum Mechanics (QM)

- Interpretation of QM ?
- What is the meaning of the qm formalism ?
- It is hard to understand in comparison to *classical physics* !

# The interpretation of quantum mechanics

This talk is about the interpretation of Quantum Mechanics (QM)

- Interpretation of QM ?
- What is the meaning of the qm formalism ?
- It is hard to understand in comparison to *classical physics* !

# The interpretation of quantum mechanics

This talk is about the interpretation of Quantum Mechanics (QM)

- Interpretation of QM ?
- What is the meaning of the qm formalism ?
- It is hard to understand in comparison to *classical physics* !

# What do we mean by *classical physics* ?

- What exactly do we mean (or should we mean) by *classical physics* ?
- Common answer: *classical mechanics* - a **deterministic** theory describing the movement of point particles.

Let us characterize the common answer by the term "particle picture".

# What do we mean by *classical physics* ?

- What exactly do we mean (or should we mean) by *classical physics* ?
- Common answer: *classical mechanics* - a **deterministic** theory describing the movement of point particles.

Let us characterize the common answer by the term "particle picture".

# Two manifestations of the particle picture.

There are two different (but possibly related) manifestations of the particle picture:

- The formal transition from classical physics to quantum mechanics, i.e. the *quantization procedure*, and
- the *interpretation* of the experimental data.

# Two manifestations of the particle picture.

There are two different (but possibly related) manifestations of the particle picture:

- The formal transition from classical physics to quantum mechanics, i.e. the *quantization procedure*, and
- the *interpretation* of the experimental data.



# The quantization procedure in the particle picture

In the particle picture the *Hamilton function*  $H(p, q)$  of classical particle physics has to be used as a starting point.

The "canonical" quantization procedure:

Replace the classical formula  $H(p, q) = E$  by  $H(\hat{p}, \hat{q})\psi = \hat{E}\psi$ ,  
where  $\hat{q} = q$ , and

$$\hat{p} = \frac{\hbar}{i} \frac{d}{dx}, \quad \hat{E} = -\frac{\hbar}{i} \frac{d}{dt}.$$

# Calling the particle picture into question

The particle picture represents the prevailing opinion. Despite of that it may be allowed to ask the following two questions:

# First question

## 1 Is it the only possible choice ?

Answer: No, QM could as well be compared to a classical *statistical* theory.

A statistical theory (as defined here) is **not deterministic** with regard to single events.

# First question

1 Is it the only possible choice ?

Answer: No, QM could as well be compared to a classical *statistical* theory.

A statistical theory (as defined here) is **not deterministic** with regard to single events.

## Second question

### 2 Is it clear and simple ?

Answer: No, certainly not

- The above quantization procedure is *incomprehensible*.
- Experimental data (interpreted in the particle picture) are *mysterious*.

## Second question

### 2 Is it clear and simple ?

Answer: No, certainly not

- The above quantization procedure is *incomprehensible*.
- Experimental data (interpreted in the particle picture) are *mysterious*.

## Second question

### 2 Is it clear and simple ?

Answer: No, certainly not

- The above quantization procedure is *incomprehensible*.
- Experimental data (interpreted in the particle picture) are *mysterious*.

## Second question

### 2 Is it clear and simple ?

Answer: No, certainly not

- The above quantization procedure is *incomprehensible*.
- Experimental data (interpreted in the particle picture) are *mysterious*.



# Where do all the mysteries come from ?

- Why does the particle picture lead to mysterious results ?

Answer: The qm formalism is *unable to make predictions about individual events*. Only probabilities are provided.

Thus, there is a **clash** between the available theory and our desire for a *complete* theory (complete means here simply: deterministic with regard to single events)

# Where do all the mysteries come from ?

- Why does the particle picture lead to mysterious results ?

**Answer:** The qm formalism is *unable to make predictions about individual events*. Only probabilities are provided.

Thus, there is a **clash** between the available theory and our desire for a *complete* theory (complete means here simply: deterministic with regard to single events)

# Where do all the mysteries come from ?

- Why does the particle picture lead to mysterious results ?

Answer: The qm formalism is *unable to make predictions about individual events*. Only probabilities are provided.

Thus, there is a **clash** between the available theory and our desire for a *complete* theory (complete means here simply: deterministic with regard to single events)

# How to escape the clash ?

There are two strategies to escape the clash between wish and reality:

- Complement QM by additional intellectual constructs, preferably of a highly abstract (in a philosophical and/or mathematical sense) character.
- Give up the particle picture – i.e. consider QM as a purely statistical theory (basically Einsteins view of QM, see Ballentines article in Reviews of Modern Physics).

# How to escape the clash ?

There are two strategies to escape the clash between wish and reality:

- Complement QM by additional intellectual constructs, preferably of a highly abstract (in a philosophical and/or mathematical sense) character.
- Give up the particle picture – i.e. consider QM as a purely statistical theory (basically Einsteins view of QM, see Ballentines article in Reviews of Modern Physics).

# A calculation supporting an interpretation

- In the rest of this talk I will try to make a point in favor of this Statistical Interpretation of QM.
- The incomprehensibility of the quantization procedure and the mysterious nature of the experimental data are different aspects of *one and the same thing*.
- I assume that the particle picture is the "wrong" starting point for quantization, and that a statistical theory should be used instead.
- A successful "simple" *quantization* (derivation of QM) from statistical postulates would also present a strong argument in favor of the Statistical *Interpretation* of QM.

# A calculation supporting an interpretation

- In the rest of this talk I will try to make a point in favor of this Statistical Interpretation of QM.
- The incomprehensibility of the quantization procedure and the mysterious nature of the experimental data are different aspects of *one and the same thing*.
- I assume that the particle picture is the "wrong" starting point for quantization, and that a statistical theory should be used instead.
- A successful "simple" *quantization* (derivation of QM) from statistical postulates would also present a strong argument in favor of the Statistical *Interpretation* of QM.

# A calculation supporting an interpretation

- In the rest of this talk I will try to make a point in favor of this Statistical Interpretation of QM.
- The incomprehensibility of the quantization procedure and the mysterious nature of the experimental data are different aspects of *one and the same thing*.
- I assume that the particle picture is the "wrong" starting point for quantization, and that a statistical theory should be used instead.
- A successful "simple" *quantization* (derivation of QM) from statistical postulates would also present a strong argument in favor of the Statistical *Interpretation* of QM.



# A calculation supporting an interpretation

- In the rest of this talk I will try to make a point in favor of this Statistical Interpretation of QM.
- The incomprehensibility of the quantization procedure and the mysterious nature of the experimental data are different aspects of *one and the same thing*.
- I assume that the particle picture is the "wrong" starting point for quantization, and that a statistical theory should be used instead.
- A successful "simple" *quantization* (derivation of QM) from statistical postulates would also present a strong argument in favor of the Statistical *Interpretation* of QM.

# Where to find more details

I will be unable to report details on the calculation in this talk.  
More information may be found

- in a preprint to be uploaded to the Arxiv e-print repository in the near future
- in parts of a draft at Arxiv, which is, however, incomplete and inconclusive,
- in a previous (precursor) paper using the same philosophy and a different set of postulates:  
U.Klein, "Schrödinger's equation with gauge coupling derived from a continuity equation", Found. Phys. **39** 964-995 (2009)

# Where to find more details

I will be unable to report details on the calculation in this talk.  
More information may be found

- in a preprint to be uploaded to the Arxiv e-print repository in the near future
- in parts of a draft at Arxiv, which is, however, incomplete and inconclusive,
- in a previous (precursor) paper using the same philosophy and a different set of postulates:  
U.Klein, "Schrödinger's equation with gauge coupling derived from a continuity equation", Found. Phys. **39** 964-995 (2009)

# Where to find more details

I will be unable to report details on the calculation in this talk.  
More information may be found

- in a preprint to be uploaded to the Arxiv e-print repository in the near future
- in parts of a draft at Arxiv, which is, however, incomplete and inconclusive,
- in a previous (precursor) paper using the same philosophy and a different set of postulates:  
U.Klein, "Schrödinger's equation with gauge coupling derived from a continuity equation", Found. Phys. **39** 964-995 (2009)

# The theory, general features

- Classical statistical physics contains both deterministic and indeterministic elements. The idea is to **remove the deterministic elements**
- Dynamical laws do only exist for ensemble averages

# The theory, general features

- Classical statistical physics contains both deterministic and indeterministic elements. The idea is to **remove the deterministic elements**
- Dynamical laws do only exist for ensemble averages

# How to obtain the fundamental dynamical laws ?

- A possible rule to obtain a fundamental condition for such "generalized" statistical theories:  
start from a particle theory and then,  
*replace in the particle theory all observables by ensemble averages.*

This leads to our "first assumption" (see next frame)

- We study a single classical particle in a single spatial dimension  $x$  .

# How to obtain the fundamental dynamical laws ?

- A possible rule to obtain a fundamental condition for such "generalized" statistical theories:  
start from a particle theory and then,  
*replace in the particle theory all observables by ensemble averages.*

This leads to our "first assumption" (see next frame)

- We study a single classical particle in a single spatial dimension  $x$  .



# The first assumption

The "*true dynamical law*" for a particle in a force field

$$F(x) = -\frac{dV(x)}{dx},$$

is not

$$\frac{d}{dt}x(t) = \frac{p(t)}{m}, \quad \frac{d}{dt}p(t) = F(x(t)),$$

but

$$\frac{d}{dt}\bar{x} = \frac{\bar{p}}{m} \quad \frac{d}{dt}\bar{p} = \overline{F(x)},$$

where  $\bar{x}$ ,  $\bar{p}$  and  $\overline{F}$  are *ensemble averages*.

# Some comments on the first assumption

- The quantities  $x$ ,  $p$  are not observables but random variables.
- The ensemble averages are given by

$$\bar{x} = \int_{-\infty}^{\infty} dx \rho(x, t) x, \quad \bar{p} = \int_{-\infty}^{\infty} dp w(p, t) p, \dots \text{etc,}$$

- with the probabilities  $\rho(x, t)$  and  $w(p, t)$  playing the role of the new "observables".
- No law is yet known for  $\rho(x, t)$  and  $w(p, t)$ . *Many different laws* are possible  $\implies$  the first assumption may be referred to as "statistical condition(s)".

# Some comments on the first assumption

- The quantities  $x$ ,  $p$  are not observables but random variables.
- The ensemble averages are given by

$$\bar{x} = \int_{-\infty}^{\infty} dx \rho(x, t) x, \quad \bar{p} = \int_{-\infty}^{\infty} dp w(p, t) p, \dots \text{etc,}$$

- with the probabilities  $\rho(x, t)$  and  $w(p, t)$  playing the role of the new "observables".
- No law is yet known for  $\rho(x, t)$  and  $w(p, t)$ . *Many different laws* are possible  $\implies$  the first assumption may be referred to as "statistical condition(s)".

# Some comments on the first assumption

- The quantities  $x$ ,  $p$  are not observables but random variables.
- The ensemble averages are given by

$$\bar{x} = \int_{-\infty}^{\infty} dx \rho(x, t) x, \quad \bar{p} = \int_{-\infty}^{\infty} dp w(p, t) p, \dots \text{etc,}$$

- with the probabilities  $\rho(x, t)$  and  $w(p, t)$  playing the role of the new "observables".
- No law is yet known for  $\rho(x, t)$  and  $w(p, t)$ . *Many different laws* are possible  $\implies$  the first assumption may be referred to as "statistical condition(s)".

# Some comments on the first assumption

- The quantities  $x$ ,  $p$  are not observables but random variables.
- The ensemble averages are given by

$$\bar{x} = \int_{-\infty}^{\infty} dx \rho(x, t) x, \quad \bar{p} = \int_{-\infty}^{\infty} dp w(p, t) p, \dots \text{etc,}$$

- with the probabilities  $\rho(x, t)$  and  $w(p, t)$  playing the role of the new "observables".
- No law is yet known for  $\rho(x, t)$  and  $w(p, t)$ . *Many different laws* are possible  $\implies$  the first assumption may be referred to as "statistical condition(s)".

# The second assumption

A conservation law for the probability density  $\rho(x, t)$ ,

$$\frac{\partial \rho(x, t)}{\partial t} + \frac{\partial}{\partial x} \frac{\rho(x, t)}{m} p(x, t) = 0,$$

exists, where the momentum field  $p(x, t)$  can be written as a gradient of a scalar function  $S(x, t)$

$$p(x, t) = \frac{\partial S(x, t)}{\partial x}.$$

The momentum field  $p(x, t)$  and the random variable  $p$  are (generally) different !

# We need a third assumption

- The above two assumptions lead to an infinite number of statistical theories. [Differential equations for  $\rho(x, t)$  and  $S(x, t)$ ].
- These theories may be labelled by an (almost) arbitrary function  $L$ . Special cases are QM (for  $L = L^q$ ) and a classical statistical theory, which is the classical limit of QM (for  $L = 0$ ).
- Why is  $L = L^q$  realized and not any other choice ? Which principle has been used by nature ?

# We need a third assumption

- The above two assumptions lead to an infinite number of statistical theories. [Differential equations for  $\rho(x, t)$  and  $S(x, t)$ ].
- These theories may be labelled by an (almost) arbitrary function  $L$ . Special cases are QM (for  $L = L^q$ ) and a classical statistical theory, which is the classical limit of QM (for  $L = 0$ ).
- Why is  $L = L^q$  realized and not any other choice ? Which principle has been used by nature ?



# We need a third assumption

- The above two assumptions lead to an infinite number of statistical theories. [Differential equations for  $\rho(x, t)$  and  $S(x, t)$ ].
- These theories may be labelled by an (almost) arbitrary function  $L$ . Special cases are QM (for  $L = L^q$ ) and a classical statistical theory, which is the classical limit of QM (for  $L = 0$ ).
- **Why is  $L = L^q$  realized and not any other choice ? Which principle has been used by nature ?**

# The third assumption

We hope that something like the principle of maximal entropy (in statistical thermodynamics) might work here, and postulate:

- The principle used by nature is the **principle of maximal disorder**, as realized by the principle of *minimal Fisher information*

# These assumptions imply Schrödinger's equation

The above three assumptions imply

$$-\frac{\hbar}{i} \frac{\partial \psi}{\partial t} = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi}{\partial x^2} + V\psi,$$

where

$$\psi = \sqrt{\rho} e^{i\frac{S}{\hbar}}.$$

- Two open questions ?
- $w(p, t)$ ?
- How to obtain expectation values of  $p$ -dependend quantities ?

# These assumptions imply Schrödinger's equation

The above three assumptions imply

$$-\frac{\hbar}{i} \frac{\partial \psi}{\partial t} = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi}{\partial x^2} + V\psi,$$

where

$$\psi = \sqrt{\rho} e^{i\frac{S}{\hbar}}.$$

- Two open questions ?
- $w(p, t)$ ?
- How to obtain expectation values of  $p$ -dependend quantities ?

# These assumptions imply Schrödinger's equation

The above three assumptions imply

$$-\frac{\hbar}{i} \frac{\partial \psi}{\partial t} = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi}{\partial x^2} + V\psi,$$

where

$$\psi = \sqrt{\rho} e^{i\frac{S}{\hbar}}.$$

- Two open questions ?
- $w(p, t)$ ?
- How to obtain expectation values of  $p$ -dependend quantities ?

# These assumptions imply Schrödinger's equation

The above three assumptions imply

$$-\frac{\hbar}{i} \frac{\partial \psi}{\partial t} = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi}{\partial x^2} + V\psi,$$

where

$$\psi = \sqrt{\rho} e^{i\frac{S}{\hbar}}.$$

- Two open questions ?
- $w(p, t)$ ?
- How to obtain expectation values of  $p$ -dependend quantities ?

# The momentum random variable

We constructed a "configuration space" probabilistic theory, i.e.:  
The sample space is  $\mathcal{R}$ , the set of possible outcomes of  $x$ .

On the other hand:

- The momentum  $p$  of the present theory is a random variable.
- But this random variable  $p$  is not defined as a function of  $x$  - like a random variable of "classical" probability theory.
- This fact presents the essential **nonclassical** feature of the present approach to QM - it means that a **deterministic element** must be **removed** from standard probability theory in order to obtain QM.
- To obtain  $w(p, t)$  and to answer the above questions one more condition must be implemented.

# The momentum random variable

We constructed a "configuration space" probabilistic theory, i.e.:  
The sample space is  $\mathcal{R}$ , the set of possible outcomes of  $x$ .

On the other hand:

- The momentum  $p$  of the present theory is a random variable.
- But this random variable  $p$  is not defined as a function of  $x$  - like a random variable of "classical" probability theory.
- This fact presents the essential **nonclassical** feature of the present approach to QM - it means that a **deterministic element** must be **removed** from standard probability theory in order to obtain QM.
- To obtain  $w(p, t)$  and to answer the above questions one more condition must be implemented.



# The momentum random variable

We constructed a "configuration space" probabilistic theory, i.e.:  
The sample space is  $\mathcal{R}$ , the set of possible outcomes of  $x$ .

On the other hand:

- The momentum  $p$  of the present theory is a random variable.
- But this random variable  $p$  is not defined as a function of  $x$  - like a random variable of "classical" probability theory.
- This fact presents the essential **nonclassical** feature of the present approach to QM - it means that a **deterministic element** must be **removed** from standard probability theory in order to obtain QM.
- To obtain  $w(p, t)$  and to answer the above questions one more condition must be implemented.

# The momentum random variable

We constructed a "configuration space" probabilistic theory, i.e.:  
The sample space is  $\mathcal{R}$ , the set of possible outcomes of  $x$ .

On the other hand:

- The momentum  $p$  of the present theory is a random variable.
- But this random variable  $p$  is not defined as a function of  $x$  - like a random variable of "classical" probability theory.
- This fact presents the essential **nonclassical** feature of the present approach to QM - it means that a **deterministic element** must be **removed** from standard probability theory in order to obtain QM.
- To obtain  $w(p, t)$  and to answer the above questions one more condition must be implemented.

# The momentum random variable

We constructed a "configuration space" probabilistic theory, i.e.:  
The sample space is  $\mathcal{R}$ , the set of possible outcomes of  $x$ .

On the other hand:

- The momentum  $p$  of the present theory is a random variable.
- But this random variable  $p$  is not defined as a function of  $x$  - like a random variable of "classical" probability theory.
- This fact presents the essential **nonclassical** feature of the present approach to QM - it means that a **deterministic element** must be **removed** from standard probability theory in order to obtain QM.
- To obtain  $w(p, t)$  and to answer the above questions one more condition must be implemented.

# The momentum random variable

We constructed a "configuration space" probabilistic theory, i.e.:  
The sample space is  $\mathcal{R}$ , the set of possible outcomes of  $x$ .

On the other hand:

- The momentum  $p$  of the present theory is a random variable.
- But this random variable  $p$  is not defined as a function of  $x$  - like a random variable of "classical" probability theory.
- This fact presents the essential **nonclassical** feature of the present approach to QM - it means that a **deterministic element** must be **removed** from standard probability theory in order to obtain QM.
- To obtain  $w(p, t)$  and to answer the above questions one more condition must be implemented.

# Energy conservation in the mean

Assumption no.4 is *energy conservation in the mean*:

$$\frac{d}{dt} \left[ \int_{-\infty}^{\infty} dp w(p, t) \frac{p^2}{2m} + \int_{-\infty}^{\infty} dx \rho(x, t) V(x) \right] = 0.$$

- Assumptions 1-4 imply:  
As far as the calculation of expectation values of  $p^n$  for  $n = 0, 1, 2$  is concerned, the probability density is given by

$$w(p, t) = \frac{1}{\hbar} |\phi(p, t)|^2,$$

where  $\phi(p, t)$  is the Fourier transform of  $\psi(x, t)$ .

- No values of  $n$  different from  $n = 0, 1, 2$  seem to exist in realistic situations.

# Energy conservation in the mean

Assumption no.4 is *energy conservation in the mean*:

$$\frac{d}{dt} \left[ \int_{-\infty}^{\infty} dp w(p, t) \frac{p^2}{2m} + \int_{-\infty}^{\infty} dx \rho(x, t) V(x) \right] = 0.$$

- Assumptions 1-4 imply:

As far as the calculation of expectation values of  $p^n$  for  $n = 0, 1, 2$  is concerned, the probability density is given by

$$w(p, t) = \frac{1}{\hbar} |\phi(p, t)|^2,$$

where  $\phi(p, t)$  is the Fourier transform of  $\psi(x, t)$ .

- No values of  $n$  different from  $n = 0, 1, 2$  seem to exist in realistic situations.

# Energy conservation in the mean

Assumption no.4 is *energy conservation in the mean*:

$$\frac{d}{dt} \left[ \int_{-\infty}^{\infty} dp w(p, t) \frac{p^2}{2m} + \int_{-\infty}^{\infty} dx \rho(x, t) V(x) \right] = 0.$$

- Assumptions 1-4 imply:

As far as the calculation of expectation values of  $p^n$  for  $n = 0, 1, 2$  is concerned, the probability density is given by

$$w(p, t) = \frac{1}{\hbar} |\phi(p, t)|^2,$$

where  $\phi(p, t)$  is the Fourier transform of  $\psi(x, t)$ .

- No values of  $n$  different from  $n = 0, 1, 2$  seem to exist in realistic situations.

# Summary

- We found a set of assumptions which imply QM. All assumptions may be interpreted in physical terms.
- This *comprehensible* quantization procedure does not start from a classical particle Hamiltonian but from a (abstract) statistical theory.
- It gives an explanation for the success of the canonical quantization procedure - for rules like  $p \rightarrow \frac{\hbar}{i} \frac{d}{dx}$



# Summary

- We found a set of assumptions which imply QM. All assumptions may be interpreted in physical terms.
- This *comprehensible* quantization procedure does not start from a classical particle Hamiltonian but from a (abstract) statistical theory.
- It gives an explanation for the success of the canonical quantization procedure - for rules like  $p \rightarrow \frac{\hbar}{i} \frac{d}{dx}$

# Summary

- We found a set of assumptions which imply QM. All assumptions may be interpreted in physical terms.
- This *comprehensible* quantization procedure does not start from a classical particle Hamiltonian but from a (abstract) statistical theory.
- It gives an explanation for the success of the canonical quantization procedure - for rules like  $p \rightarrow \frac{\hbar}{i} \frac{d}{dx}$

# Summary

- The central new assumption (no.1) was a set of two dynamical equations with the structure of classical particle mechanics - but with ensemble averages replacing observables.
- QM is a configuration space probabilistic theory with a new element of indeterminism (momentum no longer standard random variable) added.
- All this gives very strong support to the **Statistical Interpretation** of QM.

# Summary

- The central new assumption (no.1) was a set of two dynamical equations with the structure of classical particle mechanics - but with ensemble averages replacing observables.
- QM is a configuration space probabilistic theory with a new element of indeterminism (momentum no longer standard random variable) added.
- All this gives very strong support to the **Statistical Interpretation** of QM.

# Summary

- The central new assumption (no.1) was a set of two dynamical equations with the structure of classical particle mechanics - but with ensemble averages replacing observables.
- QM is a configuration space probabilistic theory with a new element of indeterminism (momentum no longer standard random variable) added.
- All this gives very strong support to the **Statistical Interpretation** of QM.